# On the product of distributions with coincident point singularities * 

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#### Abstract

In this paper we generalize the result obtained by Gonzáles Domínguez, Scarfiello and Fisher (A. Gonzáles Domínguez and R. Scarfiello, Rev. Un. Mat. Argentina, Volumen de Homenaje a Beppo Levi (1956) 58-67; B. Fisher, Proc. Cambridge Philos. Soc. 72 (1972) 201-204). This result can be used in quantum field theory for the evaluation of products of propagators of the fields. With this new result we obtain the product $\left(c^{2}+P\right)^{-n} \cdot \delta^{m-1}\left(c^{2}+P\right)$. As a physical example, we evaluate the self-energy Green function of a massless scalar field.


## 1. Introduction

Our purpose is to evaluate the product $x^{-n} \cdot \delta^{m-1}(x)(n, m=1,2, \ldots)$, which appears in applications, essencially in the quantum theory of fields, in the evaluation of products of propagators.

The distribution $\delta^{(s)}(x)$ is the $s$ th derivative of the $\delta(x)$ distribution and $x^{-n}$ ( $n=1,2, \ldots$ ) is the distribution defined by the formula (see [8])

$$
\begin{equation*}
x^{-n}=x_{+}^{-n}+(-1)^{n} x_{-}^{n} \tag{1}
\end{equation*}
$$

where (see [11, pp. 86, 89])

$$
\begin{align*}
\left(x_{+}^{-n}, \varphi\right)= & \int_{0}^{\infty} x^{-n}\left[\varphi(x)-\varphi(0)-x \varphi^{\prime}(0)-\cdots\right. \\
& \left.-\frac{x^{(n-1)}}{(n-1)!} \varphi^{(n-1)}(0) H(1-x)\right] \mathrm{d} x \tag{2}
\end{align*}
$$

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\[

$$
\begin{align*}
\left(x_{-}^{-n}, \varphi\right)= & \int_{0}^{\infty} x^{-n}\left[\varphi(-x)-\varphi(0)+x \varphi^{\prime}(0)+\cdots\right. \\
& \left.+(-1)^{n-1} \frac{x^{(n-1)}}{(n-1)!} \varphi^{(n-1)}(0) H(1-x)\right] \mathrm{d} x \tag{3}
\end{align*}
$$
\]

In (2) and (3), $H(x)$ is Heaviside's step function.
The distribution $x^{-n}$ is sometimes denoted by $f p x^{-n}$, where the letters $f p$ mean "finite part", for $n>1 ; x^{-1}$ is usually denoted by $P v x^{-1}$, where the letters $P v$ mean "principal value".

To obtain the product $x^{-n} \cdot \delta^{(m-1)}(x)$ we also need the following formulae (see [11, pp. 93, 94]):

$$
\begin{align*}
(x \pm \mathrm{i} 0)^{-n} & =x^{-n} \mp \pi \mathrm{i} \frac{(-1)^{n-1}}{(n-1)!} \delta^{(n-1)}(x)  \tag{4}\\
(x \pm \mathrm{i} 0)^{\lambda} & =x_{+}^{\lambda}+\mathrm{e}^{ \pm \lambda \pi \mathrm{i}} x_{-}^{\lambda}  \tag{5}\\
x_{+}^{\lambda} & =\frac{(-1)^{n-1} \delta^{(n-1)}(x)}{(n-1)!(\lambda+n)}+F_{-n}\left(x_{+}, \lambda\right)  \tag{6}\\
x_{-}^{\lambda} & =\frac{\delta^{(n-1)}(x)}{(n-1)!(\lambda+n)}+F_{-n}\left(x_{-}, \lambda\right) \tag{7}
\end{align*}
$$

where $F_{-n}\left(x_{+}, \lambda\right)$ and $F_{-n}\left(x_{-}, \lambda\right)$ are the regular parts of the Laurent expansions (see $[11, \mathrm{pp} .86,88]$ ) and $x_{ \pm}^{\lambda}$ are the generalized functions defined by

$$
x_{+}^{\lambda}= \begin{cases}x^{\lambda} & \text { for } x>0  \tag{8}\\ 0 & \text { for } x \leqslant 0\end{cases}
$$

and

$$
x_{-}^{\lambda}= \begin{cases}|x|^{\lambda} & \text { for } x<0  \tag{9}\\ 0 & \text { for } x \geqslant 0\end{cases}
$$

Also we know that the following formulae are valid [10]:

$$
\begin{align*}
& \Gamma(z+n)=z(z+1) \cdots(z+n-1) \Gamma(z)  \tag{10}\\
& \frac{\Gamma(z)}{\Gamma(z-n)}=(-1)^{n} \frac{\Gamma(-z+n+1)}{\Gamma(-z+1)}  \tag{11}\\
& \Gamma\left(\frac{1}{2}+z\right) \Gamma\left(\frac{1}{2}-z\right)=\pi \sec (\pi z)  \tag{12}\\
& \Gamma(2 z)=2^{2 z-1} \pi^{-1 / 2} \Gamma(z) \Gamma\left(z+\frac{1}{2}\right) \tag{13}
\end{align*}
$$

The latter is also known as the "duplication formulae" for Euler's $\Gamma$-function. In order to deal with multidimensional generalizations of all products evaluated in sections 2 and 3 we introduce the following definition:

Let $\phi(s)$ be a distribution of one variable $s$ and let $\mathcal{U}\left(x_{1}, x_{2}, \ldots, x_{n}\right) \in C^{\infty}\left(R^{n}\right)$ be such that the $(n-1)$-dimensional manifold $\mathcal{U}\left(x_{1}, x_{2}, \ldots, x_{n}\right)=0$ has no critical points; $\phi_{\mathcal{U}}(x)$ denotes the distribution defined by the formula (called the Leray formula [15])

$$
\begin{equation*}
\left\langle\phi_{\mathcal{U}}(x, \varphi(x))\right\rangle=\langle\phi(s), \psi(s)\rangle \tag{14}
\end{equation*}
$$

where

$$
\begin{equation*}
\psi(s)=\int_{\mathcal{U}(x)=s} \varphi(x) W_{\mathcal{U}}(x, \mathrm{~d} x) \tag{15}
\end{equation*}
$$

Here $W_{\mathcal{U}}$ is an $(n-1)$-dimensional form on $\mathcal{U}$ defined as follows:

$$
\mathrm{d} \mathcal{U} \wedge W_{\mathcal{U}}=\mathrm{d} x_{1} \wedge \mathrm{~d} x_{2} \wedge \cdots \wedge \mathrm{~d} x_{n}
$$

and the orientation of the manifold $\mathcal{U}\left(x_{1}, x_{2}, \ldots, x_{n}\right)=s$ is such that

$$
\operatorname{det} W_{\mathcal{U}}(x, \mathrm{~d} x)>0
$$

## 2. The multiplicative product $x^{-n} \cdot \delta^{(m-1)}(x)$

In this section we shall obtain the formula

$$
\begin{equation*}
x_{-n} \cdot \delta^{(m-1)}(x)=\frac{(-1)^{n}(m-1)!}{2(m+n-1)!} \delta^{(m+n-1)}(x), \quad n, m \in \mathcal{N}, n, m \geqslant 1 \tag{16}
\end{equation*}
$$

Here $\mathcal{N}=$ natural numbers.
We shall study this product for $n$ and $m$ positive integers taking into account the following three cases:
(1) $m=n$,
(2) $m>n$,
(3) $m<n$.

Formula (16) has been proved for $m=n$ by Gonzáles Domínguez and Scarfiello in [9]. It was later rediscovered by other authors (cf. [7,14]).

For $m=n=1$ we obtain the well-known result [9]

$$
\begin{equation*}
x^{-1} \cdot \delta(x)=-\frac{1}{2} \delta^{\prime}(x) \tag{17}
\end{equation*}
$$

According to [11, pp. 96, 97] we can write

$$
\begin{aligned}
x_{+}^{\lambda}+\mathrm{e}^{\mathrm{i} \pi \lambda} x_{-}^{\lambda}= & {\left[x^{-n}+\frac{\pi \mathrm{i}(-1)^{n}}{(n-1)!} \delta^{(n-1)}(x)\right] } \\
& +(\lambda+n)\left[\pi \mathrm{i}(-1)^{n} x_{-}^{-n}+(-1)^{n-1} \frac{\pi^{2}}{2} \frac{\delta^{(n-1)(x)}}{(n-1)!}+x^{-n} \ln |x|\right]
\end{aligned}
$$

$$
\begin{gather*}
+\frac{(\lambda+n)^{2}}{2}\left[(-1)^{n-1} \mathrm{i} \frac{\pi^{2}}{3} \frac{\delta^{(n-1)(x)}}{(n-1)!}+(-1)^{n-1} \pi^{2} x_{-}^{-n}\right. \\
\left.+2 \mathrm{i} \pi(-1)^{n} x_{-}^{-n} \ln x_{-}+x^{-n} \ln ^{2}|x|\right] \tag{18}
\end{gather*}
$$

Therefore from (18) we have

$$
\begin{equation*}
\lim _{\lambda \rightarrow-n}(x+\mathrm{i} 0)^{\lambda}=\lim _{\lambda \rightarrow-n}\left(x_{+}^{\lambda}+\mathrm{e}^{\mathrm{i} \pi \lambda} x_{-}^{\lambda}\right)=x^{-n}+\frac{\pi \mathrm{i}(-1)^{n}}{(n-1)!} \delta^{(n-1)}(x) \tag{19}
\end{equation*}
$$

From (19) and (2) we conclude that

$$
\begin{equation*}
\lim _{\lambda \rightarrow-n}(x \pm \mathrm{i} 0)^{\lambda}=(x \pm \mathrm{i} 0)^{-n} \tag{20}
\end{equation*}
$$

On the other hand, according to [11, p. 57],

$$
\begin{equation*}
\lim _{\lambda \rightarrow-n} \frac{x_{+}^{\lambda}}{\Gamma(\lambda+1)}=\delta^{(n-1)}(x), \quad n=1,2, \ldots \tag{21}
\end{equation*}
$$

where $x_{+}^{\lambda}$ is given by (8).
Case 1: The product $x^{-n} \cdot \delta^{(n-1)}(x)$
The product $x^{-n} \cdot \delta^{(n-1)}(x)$ was presented in [9] and, more recently, in [7].
We now obtain (16) for $m=n$ using (19), (13) and the formulae of [4,5,16]. From (4), (5) and considering (20) we have

$$
\begin{equation*}
x^{-n}=\frac{1}{2}\left((x+\mathrm{i} 0)^{-n}+(x-\mathrm{i} 0)^{-n}\right)=\lim _{\lambda \rightarrow-n} \frac{1}{2}\left((x+\mathrm{i} 0)^{\lambda}+(x-\mathrm{i} 0)^{\lambda}\right) \tag{22}
\end{equation*}
$$

Also, the following formulae are valid (see [5]):

$$
\begin{equation*}
x_{ \pm}^{\lambda} \cdot x_{ \pm}^{\mu}=x_{ \pm}^{\lambda+\mu} \tag{23}
\end{equation*}
$$

and

$$
\begin{equation*}
x_{+}^{\lambda} \cdot x_{-}^{\mu}=0 \tag{24}
\end{equation*}
$$

where $\lambda$ and $\mu$ are complex numbers such that $\lambda, \mu$ and $\lambda+\mu \neq-k, k=1,2, \ldots$. Formulae (21) and (22) allow us to represent $\delta^{(m)}$ and $x^{-n}$ as "canonically regularized" by these equations. In this way the canonically regularized form of the product we are looking for can be defined as the product of the corresponding $\lambda$-dependent expressions:

$$
\begin{align*}
x^{-n} \cdot \delta^{(n-1)}(x) & =\lim _{\lambda \rightarrow-n} \frac{1}{2}\left((x+\mathrm{i} 0)^{\lambda}+(x-\mathrm{i} 0)^{\lambda}\right) \cdot \frac{x_{+}^{\lambda}}{\Gamma(\lambda+1)} \\
& =\lim _{\lambda \rightarrow-n}\left[x_{+}^{\lambda}+\cos \pi \lambda x_{-}^{\lambda}\right] \cdot \frac{x_{+}^{\lambda}}{\Gamma(\lambda+1)} \\
& =\lim _{\lambda \rightarrow-n}\left[\frac{x_{+}^{\lambda} x_{+}^{\lambda}}{\Gamma(\lambda+1)}+\cos \pi \lambda \frac{x_{-}^{\lambda} x_{+}^{\lambda}}{\Gamma(\lambda+1)}\right] \tag{25}
\end{align*}
$$

From (23)-(25) we get

$$
\begin{equation*}
x^{-n} \cdot \delta^{(n-1)}(x)=\lim _{\lambda \rightarrow-n} \frac{x_{+}^{2 \lambda}}{\Gamma(\lambda+1)} . \tag{26}
\end{equation*}
$$

By using the duplication formula (13) and taking the limit, we obtain

$$
\begin{equation*}
x^{-n} \cdot \delta^{(n-1)}(x)=\frac{(-1)^{n}}{2} \frac{(n-1)!}{(2 n-1)!} \delta^{(2 n-1)}(x) . \tag{27}
\end{equation*}
$$

Formula (27) coincides with the result of $[7,9]$.
Case 2: The product $x^{-n} \cdot \delta^{(n-1)}(x)$ for $m>n$
If $m>n$, there exists $p \in \mathcal{N}$ such that $m=n+p$. Therefore

$$
\begin{equation*}
\lim _{\mu \rightarrow-m} \frac{x_{+}^{\mu}}{\Gamma(\mu+1)}=\lim _{\gamma \rightarrow-n} \frac{x_{+}^{\gamma-p}}{\Gamma(\gamma+1-p)} . \tag{28}
\end{equation*}
$$

Then, from (21), (22) and (28) we have

$$
\begin{align*}
x^{-n} \cdot \delta^{(m-1)}(x) & =\lim _{\gamma \rightarrow-n} \frac{1}{2}\left[(x+\mathrm{i} 0)^{\gamma}+(x-\mathrm{i} 0)^{\gamma}\right] \frac{x_{+}^{\gamma-p}}{\Gamma(\gamma+1-p)} \\
& =\lim _{\gamma \rightarrow-n}\left[x_{+}^{\gamma}+\cos \pi \gamma x_{-}^{\gamma}\right] \frac{x_{+}^{\gamma-p}}{\Gamma(\gamma+1-p)} . \tag{29}
\end{align*}
$$

From (23), (24), (28) and (29) we conclude that

$$
\begin{equation*}
x^{-n} \cdot \delta^{(m-1)}(x)=\lim _{\gamma \rightarrow-n} \frac{x_{+}^{2 \gamma-p}}{\Gamma(\gamma+1-p)} \tag{30}
\end{equation*}
$$

and, by taking the limit, we conclude that

$$
\begin{equation*}
x^{-n} \cdot \delta^{(m-1)}(x)=\frac{(-1)^{n}}{2} \frac{(m-1)!}{(m+n-1)!} \delta^{(m+n-1)}(x), \quad m>n . \tag{31}
\end{equation*}
$$

Case 3: The product $x^{-n} \cdot \delta^{(n-1)}(x)$ for $m<n$
If $m<n$, there exists $p \in \mathcal{N}$ such that $n=m+p$. Then, from (39) and taking into account that

$$
\begin{equation*}
\lim _{\mu \rightarrow-m} \frac{x_{+}^{\mu}}{\Gamma(\mu+1)}=\lim _{\gamma \rightarrow-n} \frac{x_{+}^{\gamma+p}}{\Gamma(\gamma+1+p)}, \tag{32}
\end{equation*}
$$

we have

$$
\begin{equation*}
x^{-n} \cdot \delta^{(m-1)}(x)=\lim _{\gamma \rightarrow-n}\left[x_{+}^{\gamma}+\cos \pi \gamma x_{-}^{\gamma}\right] \frac{x_{+}^{\gamma+p}}{\Gamma(\gamma+1+p)} . \tag{33}
\end{equation*}
$$

From (23) and (24) we obtain for (33)

$$
\begin{equation*}
x^{-n} \cdot \delta^{(m-1)}(x)=\lim _{\gamma \rightarrow-n} \frac{x_{+}^{2 \gamma+p}}{\Gamma(\gamma+1+p)} . \tag{34}
\end{equation*}
$$

And, by taking the limit, we finally get

$$
\begin{equation*}
x^{-n} \cdot \delta^{(m-1)}(x)=\frac{(-1)^{n}}{2} \frac{(m-1)!}{(m+n-1)!} \delta^{(m+n-1)}(x) \tag{35}
\end{equation*}
$$

for all $m, n \geqslant 1$. This proves our assertion (16). It is clear that with a similar procedure we can prove that

$$
\begin{equation*}
x^{-n} \cdot x^{-m}=x^{-n-m} \tag{36}
\end{equation*}
$$

and also that

$$
\begin{equation*}
\delta^{(m)}(x) \cdot \delta^{(n)}(x)=0 \tag{37}
\end{equation*}
$$

In fact, the last equation can be written as (cf. (28), $m=n+p, p \geqslant 0$ )

$$
\lim _{\lambda \rightarrow-n} \frac{x_{+}^{\lambda-p}}{\Gamma(\lambda+1-p)} \cdot \frac{x_{+}^{\lambda}}{\Gamma(\lambda+1)}=\lim _{\lambda \rightarrow-n} \frac{x_{+}^{2 \lambda-p}}{\Gamma(\lambda+1-p) \Gamma(\lambda+n)}=0
$$

The numerator has a simple pole, while the denominator has a double pole.
The last equation leads to the following theorem [4,5]:
Theorem. The product of two distributions with point support is zero.

## 3. The multiplicative product $\left(c^{2}+P\right)^{-n} \cdot \delta^{(m-1)}(P)$

In this section we want to extend result (54) to certain kinds of $n$-dimensional distributions.

We will start with some definitions.
Let $x=\left(x_{1}, x_{2}, \ldots, x_{n}\right)$ be a point of the $n$-dimensional Euclidean space $\mathcal{R}^{n}$. Consider a non-degenerate quadratic form in $n$ variables of the form

$$
\begin{equation*}
P \equiv P(x)=x_{1}^{2}+x_{2}^{2}+\cdots+x_{\mu}^{2}-x_{\mu+1}^{2}-x_{\mu+2}^{2}-\cdots-x_{\mu+\nu}^{2} \tag{38}
\end{equation*}
$$

where $\mu+\nu=n=$ dimension of the space. The distributions $\left(c^{2}+P \pm \mathrm{i} 0\right)^{\lambda}$ are defined by (see [11, p. 289])

$$
\begin{equation*}
\left(c^{2}+P \pm \mathrm{i} 0\right)^{\lambda}=\lim _{\varepsilon \rightarrow 0}\left(c^{2}+P \pm \mathbf{i} \varepsilon|x|^{2}\right)^{\lambda} \tag{39}
\end{equation*}
$$

where $|x|^{2}=x_{1}^{2}+x_{2}^{2}+\cdots+x_{n}^{2}, \varepsilon>0, \lambda$ is a complex number and $c^{2}$ is a positive real number.

It is useful to state an equivalent definition of the distributions $\left(c^{2}+P \pm \mathrm{i} 0\right)^{\lambda}$. This definition involves the distributions

$$
\left(c^{2}+P\right)_{+}^{\lambda}= \begin{cases}\left(c^{2}+P\right)^{\lambda} & \text { for } c^{2}+P>0,  \tag{40}\\ 0 & \text { for } c^{2}+P \leqslant 0,\end{cases}
$$

and

$$
\left(c^{2}+P\right)_{-}^{\lambda}= \begin{cases}\left(-c^{2}-P\right)^{\lambda} & \text { for } c^{2}+P<0  \tag{41}\\ 0 & \text { for } c^{2}+P \geqslant 0\end{cases}
$$

From [6] we have

$$
\begin{equation*}
\left(c^{2}+P \pm \mathrm{i} 0\right)^{\lambda}=\left(c^{2}+P\right)_{+}^{\lambda}+\mathrm{e}^{ \pm \mathrm{i} \pi \lambda}\left(c^{2}+P\right)_{-}^{\lambda}, \tag{42}
\end{equation*}
$$

and from this formula we conclude inmediately that

$$
\begin{equation*}
\left(c^{2}+P \pm \mathrm{i} 0\right)^{\lambda}=\left(c^{2}+P\right)^{\lambda} \tag{43}
\end{equation*}
$$

when $\lambda$ is a positive integer.
It can be proved [6, p. 573, formula (2.14), p. 575, formula (3.5)] that

$$
\begin{equation*}
\left(c^{2}+P \pm \mathrm{i} 0\right)^{-k}=\left(c^{2}+P\right)^{-k} \mp \frac{(-1)^{k-1} \pi \mathrm{i}}{(k-1)!} \delta^{(k-1)}\left(c^{2}+P\right), \tag{44}
\end{equation*}
$$

where $\left(c^{2}+P\right)^{-k}$ is the regular part of the Laurent expansion of $\left(c^{2}+P\right)_{+}^{\lambda}$ near $\lambda=-k$, namely,

$$
\begin{equation*}
\left(c^{2}+P\right)^{-k}=\lim _{\lambda \rightarrow-k} \frac{\mathrm{~d}}{\mathrm{~d} \lambda}\left[(\lambda+k)\left(c^{2}+P\right)_{+}^{\lambda}\right] \tag{45}
\end{equation*}
$$

and the distribution $\delta^{(k)}\left(c^{2}+P\right)$ is defined as follows:

$$
\begin{equation*}
\left\langle\delta^{(k)}\left(c^{2}+P\right), \varphi(x)\right\rangle=\left\langle\delta^{(k)}(s), \psi(s)\right\rangle=(-1)^{k} \psi^{(k)}(0) . \tag{46}
\end{equation*}
$$

According to (15), $\psi(s)$ is given by the formula

$$
\begin{equation*}
\psi(s)=\int_{c^{2}+P=s} \varphi(x) W_{c^{2}+P}(x, \mathrm{~d} x) \tag{47}
\end{equation*}
$$

The distribution $\left(c^{2}+P\right)^{-k}$ can be defined, using (47), in the following way:

$$
\begin{align*}
& \left\langle\left(c^{2}+P\right)^{-k}, \varphi(x)\right\rangle=\left\langle s^{-k}, \psi(s)\right\rangle \\
& \quad=\int_{0}^{\infty} s^{-k}\left[\psi(s)-\psi(0)-\cdots-\frac{s^{k-1}}{(k-1)!} \psi^{(k-1)}(0) H(1-s)\right] \mathrm{d} s \tag{48}
\end{align*}
$$

According to these definitions, we can simply substitute $x$ for $c^{2}+P$ in both sides of (35) to get

$$
\begin{equation*}
\left(c^{2}+P\right)^{-n} \delta^{(m-1)}\left(c^{2}+P\right)=\frac{(-1)^{n}}{2} \frac{(m-1)!}{(m+n-1)!} \delta^{(m+n-1)}\left(c^{2}+P\right) . \tag{49}
\end{equation*}
$$

## 4. A physical example

We give now a physical example. We consider a massless scalar $(\lambda / 4!) \phi^{4}(x)$ theory in four dimensions. For this theory we shall evaluate the self-energy Green function.

The Lagrangian of this theory is

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi+\frac{\lambda}{4!} \phi^{4} . \tag{50}
\end{equation*}
$$

In the usual metric $(1,-1,-1,-1)$ the propagator for the field $\phi(x)$ is [13]

$$
\begin{equation*}
\Delta(x)=\left[-4 \pi^{2}\left(u^{2}-\mathrm{i} 0\right)\right]^{-1} \tag{51}
\end{equation*}
$$

According to equations (A.1)-(A.6) of the appendix, we can write

$$
\begin{align*}
\left(u^{2}-\mathrm{i} 0\right)^{-1}= & \left(2 x_{0}\right)^{-1}\left[\left(x_{0}-r\right)^{-1}+\left(x_{0}+r\right)^{-1}\right]+(2 r)^{-1}\left[\delta\left(x_{0}-r\right)+\delta\left(x_{0}+r\right)\right] \\
& +C \delta\left(x_{0}-r\right) \delta\left(x_{0}+r\right) \tag{52}
\end{align*}
$$

where $C$ is an arbitrary constant appearing in the definition of some distributions [13, 8.8, 8.9] (see also the appendix).

Using equation (37) and the relation given in [1], namely,

$$
\begin{equation*}
\frac{\delta\left(x_{0}-r\right) \delta\left(x_{0}+r\right)}{2 r^{2}}=\frac{\pi}{2} \delta\left(x_{0}, x_{1}, x_{2}, x_{3}\right) \tag{53}
\end{equation*}
$$

it is easy to show with the aid of (35) that

$$
\left(u^{2}-\mathrm{i} 0\right)^{-1}\left(u^{2}-\mathrm{i} 0\right)^{-1}=\left(u^{2}-\mathrm{i} 0\right)^{-2}
$$

Then, we have for the self-energy:

$$
\begin{equation*}
\Sigma(x)=(\Delta(x))^{2}=\frac{1}{16 \pi^{4}}\left(u^{2}-\mathrm{i} 0\right)^{-2}, \tag{54}
\end{equation*}
$$

where $\left(u^{2}-\mathrm{i} 0\right)^{-2}$ is defined in [13, 8.8, 8.9] (see also the appendix).

## 5. Discussion

When we use the perturbative development in quantum field theory, we have to deal with products of distributions in configuration space, or else, with convolutions in the Fourier transformed $p$-space. Unfortunately, products or convolutions (of distributions) are in general ill-defined quantities. However, in physical applications one introduces some "regularization" scheme (for example, the dimensional regularization method [2,3,12]), which allows us to give sense to divergent integrals. A similar procedure to the use of "regulators" is the canonical product of Guelfand and Shilov.

In this work we have evaluated the canonical products $x^{-n} \cdot \delta^{(m-1)}(x), x^{-n}$. $x^{-m}, \delta^{(n)}(x) \cdot \delta^{(m)}(x)$, in the sense of Guelfand-Shilov and their generalization to the $n$-dimensional space. With the use of these products and the results of $[1,4,5]$, we have
showed that it is possible to obviate the use of any regularization method to calculate (directly) the product of propagators of particles.

## Appendix. Definitions

From [13] we have

$$
\begin{align*}
\delta^{(m)}\left(u^{2}\right)= & \delta^{(m)}\left(x^{0}+r\right)\left(x^{0}-r\right)^{-m-1} \operatorname{sgn}\left(x^{0}-r\right) \\
& +\delta^{(m)}\left(x^{0}-r\right)\left(x^{0}+r\right)^{-m-1} \operatorname{sgn}\left(x^{0}+r\right) \tag{A.1}
\end{align*}
$$

where

$$
\begin{align*}
u^{2} & =x_{0}^{2}-x_{1}^{2}-\cdots-x_{n-1}^{2}  \tag{A.2}\\
r^{2} & =x_{1}^{2}+x_{2}^{2}+\cdots+x_{n-1}^{2}  \tag{A.3}\\
\left(u^{2} \pm \mathrm{i} 0\right)^{-m} & =u^{-2 m} \pm \frac{(-1)^{m}}{(m-1)!} \mathrm{i} \pi \delta^{(m-1)}\left(u^{2}\right)  \tag{A.4}\\
x^{-m} \operatorname{sgn}(x) & =\frac{(-1)^{m-1}}{(m-1)!}\left\{|x|^{-1}\right\}^{(m-1)},  \tag{A.5}\\
|x|^{-1} & =\{\operatorname{sgn}(x) \ln |x|\}^{\prime}+C \delta(x) \tag{A.6}
\end{align*}
$$

where $C$ is an arbitrary constant.

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